

Exclusive photoproduction of Υ : from HERA to Tevatron

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In collaborations with [Wolfgang Schäfer](#) and [Antoni Szczurek](#), hep-ph arXiv:0805.0717v1

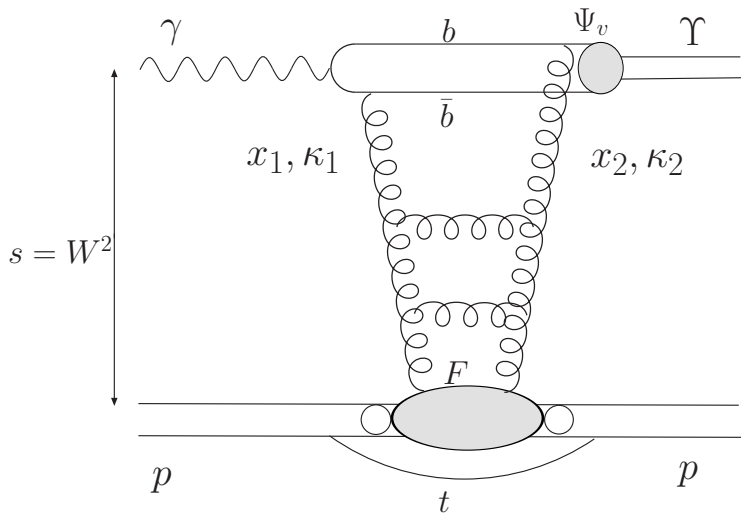
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- Exclusive photoproduction in $p\bar{p}$ collisions
 - Formalism
 - Results
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Introduction

- Exclusive production of heavy $Q\bar{Q}$ vector quarkonium states in hadronic interactions was never measured
- We restrict only to **photon - Pomeron** fusion mechanism
- The current experimental analyses at the Tevatron call for an calculation of differential distributions including the effect of absorptive corrections
- The HERA data cover energy range $W \sim 100 - 200$ GeV
- This energy range is relevant to the exclusive production at Tevatron energies for not too large rapidities of the meson

Diagram for exclusive photoproduction $\gamma p \longrightarrow \Upsilon p$



Imaginary part of the production amplitude for $\gamma p \longrightarrow \Upsilon p$

$$\Im m \mathcal{M}(W, \Delta^2 = 0, Q^2 = 0) = W^2 \frac{e\Upsilon \sqrt{4\pi\alpha_{em}}}{4\pi^2} 2 \int_0^1 \frac{dz}{z(1-z)} \int_0^\infty \pi dk^2 \psi_V(z, k^2) \\ \int_0^\infty \frac{\pi d\kappa^2}{\kappa^4} \alpha_S(q^2) \mathcal{F}(x_{eff}, \kappa^2) \left(A_0(z, k^2) W_0(k^2, \kappa^2) + A_1(z, k^2) W_1(k^2, \kappa^2) \right),$$

where

$$\begin{aligned} A_0(z, k^2) &= m_b^2 + \frac{k^2 m_b}{M + 2m_b} \\ A_1(z, k^2) &= \left[z^2 + (1-z)^2 - (2z-1)^2 \frac{m_b}{M + 2m_b} \right] \frac{k^2}{k^2 + m_b^2}, \\ W_0(k^2, \kappa^2) &= \frac{1}{k^2 + m_b^2} - \frac{1}{\sqrt{(k^2 - m_b^2 - \kappa^2)^2 + 4m_b^2 k^2}} \\ W_1(k^2, \kappa^2) &= 1 - \frac{k^2 + m_b^2}{2k^2} \left(1 + \frac{k^2 - m_b^2 - \kappa^2}{\sqrt{(k^2 - m_b^2 - \kappa^2)^2 + 4m_b^2 k^2}} \right). \end{aligned}$$

We choose the scale: $q^2 = \max\{\kappa^2, k^2 + m_b^2\}$.

Total cross section for $\gamma p \longrightarrow \Upsilon p$

The full amplitude:

$$\mathcal{M}(W, \Delta^2) = (i + \rho) \Im m \mathcal{M}(W, \Delta^2 = 0) \exp(-B(W) \Delta^2).$$

where

$$\rho = \frac{\Re e \mathcal{M}}{\Im m \mathcal{M}} = \frac{\pi}{2} \frac{\partial \log (\Im m \mathcal{M} / W^2)}{\partial \log W^2} = \frac{\pi}{2} \Delta_{\mathbf{P}}.$$

$$B(W) = B_0 + 2\alpha'_{eff} \log \left(\frac{W^2}{W_0^2} \right),$$

with: $\alpha'_{eff} = 0.164 \text{ GeV}^{-2}$, $W_0 = 95 \text{ GeV}$ (H1 Collaboration 2006)

Total cross section can be write as:

$$\sigma_{tot}(\gamma p \rightarrow \Upsilon p) = \frac{1 + \rho^2}{16\pi B(W)} \left| \Im m \frac{\mathcal{M}(W, \Delta^2)}{W^2} \right|^2$$

Parameters of the Υ wave function

Υ decay electronic width:

$$\Gamma(V \rightarrow e^+ e^-) = \frac{4\pi\alpha_{em}^2 c_\Upsilon}{3M_V^3} \cdot g_V^2 \cdot K_{NLO}$$

Leading order approximation:

$$K_{NLO} = 1$$

Next to leading order approximation:

$$K_{NLO} = 1 - \frac{16}{3\pi}\alpha_S(m_b^2)$$

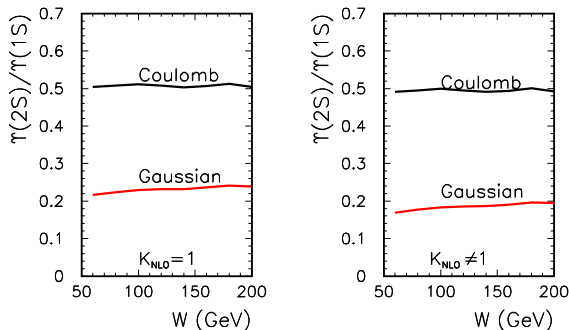
g_V -leptonic decay constant:

$$g_V = \frac{8N_c}{3} \int \frac{d^3\vec{p}}{(2\pi)^3} (M + m_b) \psi_V(p^2)$$

1) $\Gamma(V \rightarrow e^+ e^-) \Rightarrow g_V$ and this depends on K_{NLO}

2) $g_V \Rightarrow$ parameters of $\psi_V(p^2)$

The ratio of the cross section for $\Upsilon(2S)/\Upsilon(1S)$



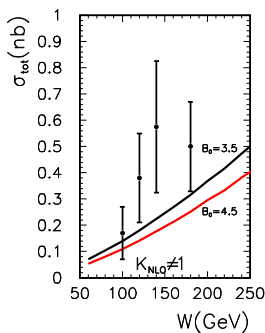
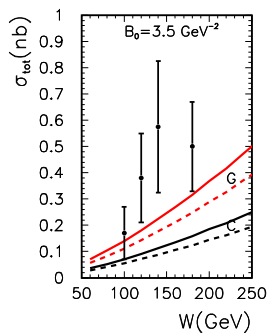
Gaussian wave function:

$$\psi_{1S}(p^2) = C_1 \exp\left(-\frac{p^2 a_1^2}{2}\right), \quad \psi_{2S}(p^2) = C_2(\xi_0 - p^2 a_2^2) \exp\left(-\frac{p^2 a_2^2}{2}\right),$$

Coulomb wave function:

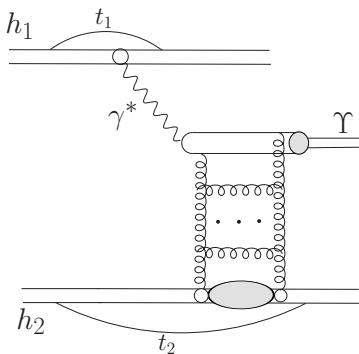
$$\psi_{1S}(p^2) = \frac{C_1}{\sqrt{M}} \frac{1}{(1 + a_1^2 p^2)^2}, \quad \psi_{2S}(p^2) = \frac{C_2}{\sqrt{M}} \frac{\xi_0 - a_2^2 p^2}{(1 + a_2^2 p^2)^3}.$$

Total cross section for the $\gamma p \rightarrow \Upsilon p$

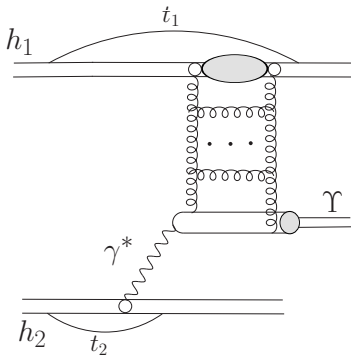


- Ψ_V — Gaussian, $K_{NLO} \neq 1$
- - - Ψ_V — Gaussian, $K_{NLO} = 1$
- Ψ_V — Coulomb, $K_{NLO} \neq 1$
- - - Ψ_V — Coulomb, $K_{NLO} = 1$

Dominant bare mechanism for $p\bar{p} \longrightarrow p\bar{p} \gamma$

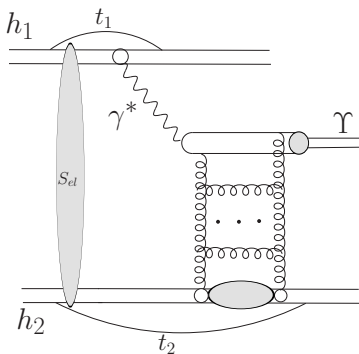


photon-pomeron

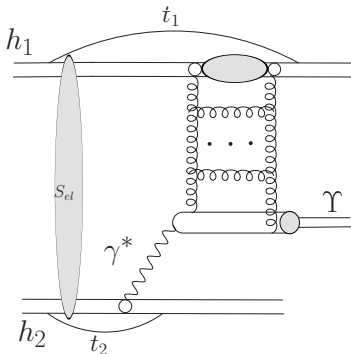


pomeron-photon

Diagram for $p\bar{p} \rightarrow p\bar{p} \Upsilon$ with absorptive correction



photon-pomeron



pomeron-photon

Amplitude for $p\bar{p} \longrightarrow p\bar{p} \Upsilon$

Amplitude without absorption:

$$\begin{aligned} M^{(0)}(\mathbf{p}_1, \mathbf{p}_2) &= e_1 \frac{2}{z_1} \frac{\mathbf{p}_1}{t_1} \mathcal{F}_{\lambda'_1 \lambda_1}(\mathbf{p}_1, t_1) \mathcal{M}_{\gamma^* h_2 \rightarrow \Upsilon h_2}(s_2, t_2, Q_1^2) \\ &+ e_2 \frac{2}{z_2} \frac{\mathbf{p}_2}{t_2} \mathcal{F}_{\lambda'_2 \lambda_2}(\mathbf{p}_2, t_2) \mathcal{M}_{\gamma^* h_1 \rightarrow \Upsilon h_1}(s_1, t_1, Q_2^2), \end{aligned}$$

Full amplitude for $p\bar{p} \longrightarrow p\bar{p} \Upsilon$:

$$\begin{aligned} M(\mathbf{p}_1, \mathbf{p}_2) &= \int \frac{d^2 \mathbf{k}}{(2\pi)^2} S_{el}(\mathbf{k}) M^{(0)}(\mathbf{p}_1 - \mathbf{k}, \mathbf{p}_2 + \mathbf{k}) \\ &= M^{(0)}(\mathbf{p}_1, \mathbf{p}_2) - \delta M(\mathbf{p}_1, \mathbf{p}_2), \end{aligned}$$

where

$$S_{el}(\mathbf{k}) = (2\pi)^2 \delta^{(2)}(\mathbf{k}) - \frac{1}{2} T(\mathbf{k}), \quad T(\mathbf{k}) = \sigma_{tot}^{pp}(s) \exp\left(-\frac{1}{2} B_{el} \mathbf{k}^2\right),$$

with $B_{el} = 17 \text{ GeV}^{-2}$, $\sigma_{tot}^{pp}(s) = 76 \text{ mb}$ (CDF Collaboration 1994)

Cross section for exclusive photoproduction in $p\bar{p}$ collisions

The absorptive correction for amplitude :

$$\delta \mathbf{M}(\mathbf{p}_1, \mathbf{p}_2) = \int \frac{d^2 \mathbf{k}}{2(2\pi)^2} T(\mathbf{k}) \mathbf{M}^{(0)}(\mathbf{p}_1 - \mathbf{k}, \mathbf{p}_2 + \mathbf{k}).$$

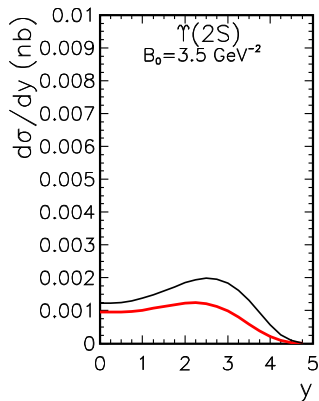
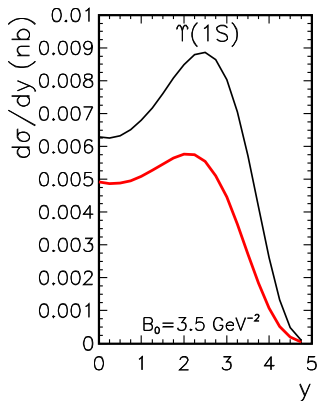
The differential cross section is given in terms of $\mathbf{M}$ as

$$d\sigma = \frac{1}{512\pi^4 s^2} |\mathbf{M}|^2 dy dt_1 dt_2 d\phi,$$

where:

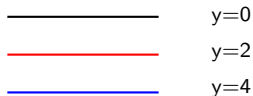
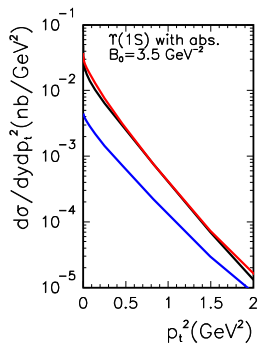
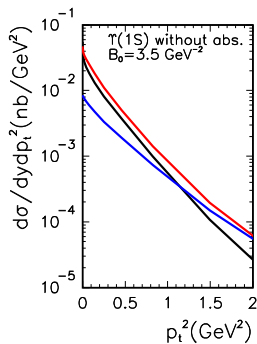
- y is rapidity of the vector meson
- $t_1 \approx -\mathbf{p}_1^2$, $t_2 \approx -\mathbf{p}_2^2$
- ϕ is azimuthal angle between $\mathbf{p}_1$ and $\mathbf{p}_2$

$d\sigma/dy$ for $\Upsilon(1S)$ and $\Upsilon(2S)$, $W = 1960$ GeV

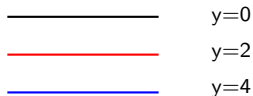
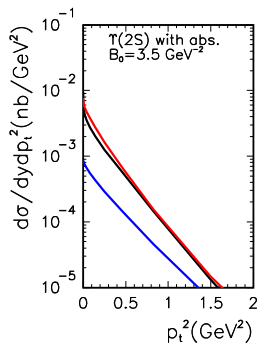
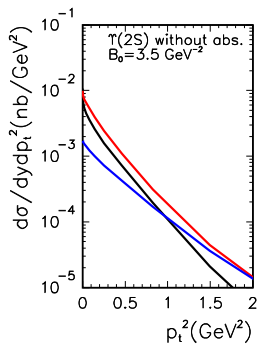


— with absorption
— without absorption

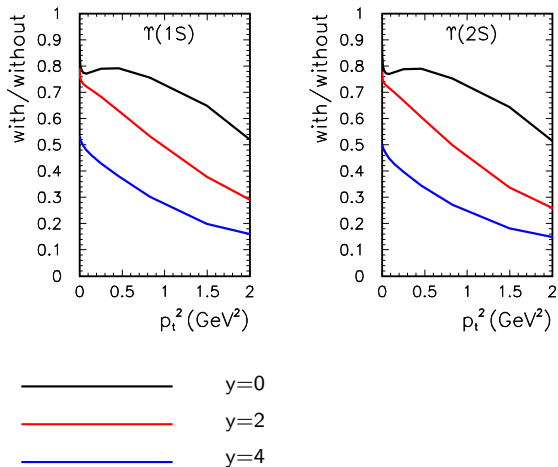
$d\sigma/dydp_t^2$ for as a function of p_t^2 for $\Upsilon(1S)$



$d\sigma/dydp_t^2$ for as a function of p_t^2 for $\Upsilon(2S)$

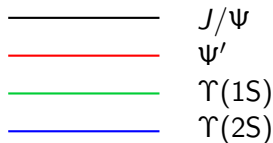
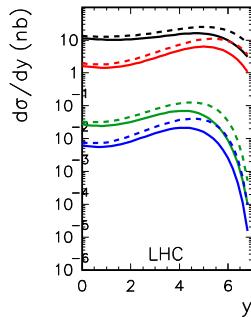
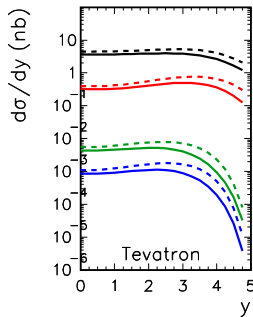


Absorption effects



The bigger rapidity $\longrightarrow$ the bigger absorption effect

$d\sigma/dy - J/\psi, \psi', \Upsilon(1S)$ and $\Upsilon(2S)$ for Tevatron and LHC energy



Conclusions

- The results for $\gamma p \rightarrow \Upsilon(1S) p$ production depend on the model of the wave function
- We have compared our results with a recent HERA data
- Our results are somewhat lower than the HERA data
- We have made predictions for $p\bar{p} \rightarrow p\bar{p} \Upsilon$ for the Tevatron (and LHC)
- Absorptive corrections are included. Their effect depends not only total cross section, but also lead to distortions of p_t and y distributions