

# Precise determination of the $f_0(600)$ and $f_0(980)$ parameters by fitting the data and dispersion relations

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- **Theory** - dispersion relations
- **Experiment** - old and present state of art
- **Results** - combined analysis of dispersion relations and data
- **Conclusions**

# Dispersion relations for $\pi\pi$ interactions theory $\longleftrightarrow$ experiment

Once subtracted dispersion relations "GKPY" (for the S and P waves):

crossing symmetry:

$$\rightarrow \vec{T}_s(s, t) = \hat{C}_{st} \vec{T}_t(t, s)$$

$\vec{T}(s, t)$  + crossing symmetry  $\rightarrow$  dispersion relations for  $4m_\pi^2 < s < \sim (1150 \text{ MeV})^2$

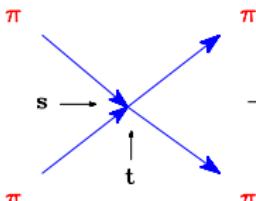
$$\text{Re } t_\ell^{I(\text{OUT})}(s) = \sum_{l'=0}^2 C_{st}^{ll'} a_0^{l'} + \sum_{l'=0}^2 \sum_{\ell'=0}^4 \int_{4m_\pi^2}^\infty ds' K_{\ell\ell'}^{ll'}(s, s') \text{Im } t_{\ell'}^{I(\text{IN})}(s')$$

$a_0^{l'}$  - subtraction constant  $= \vec{T}_s(s = 4m_\pi^2, t = 0)$  - scattering lengths in S wave

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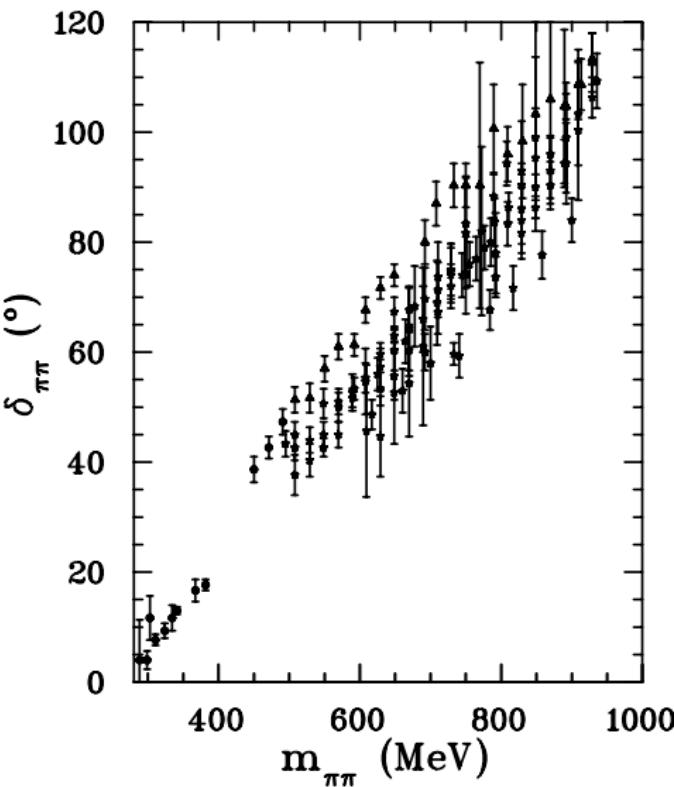
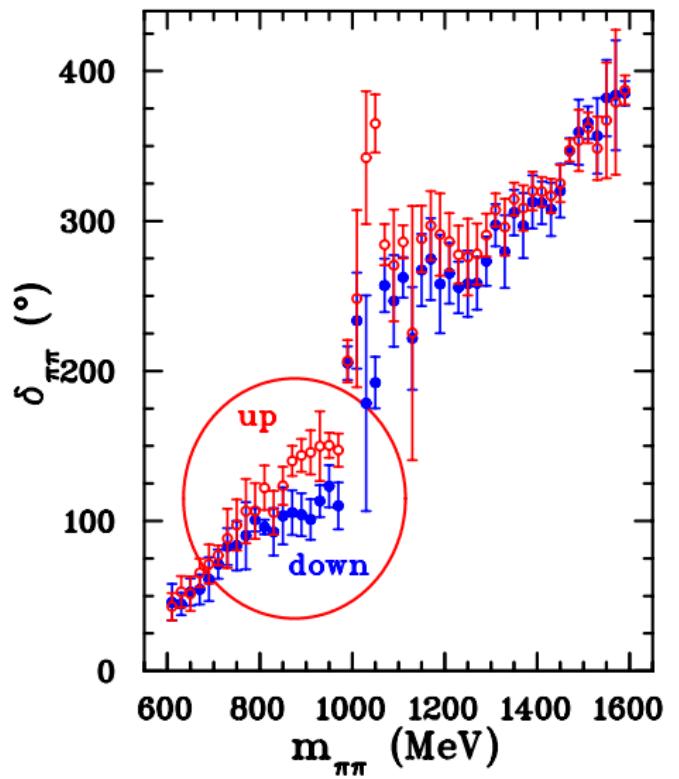
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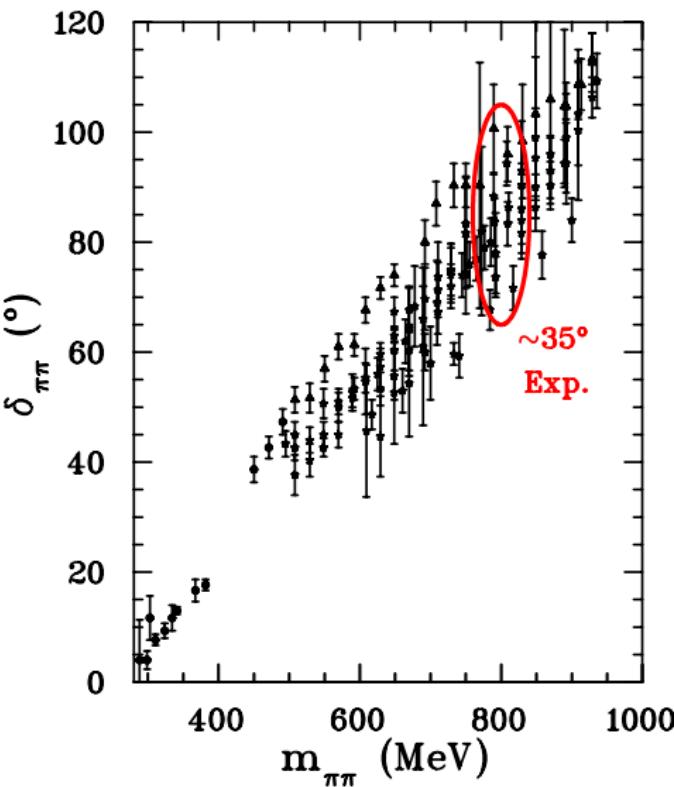
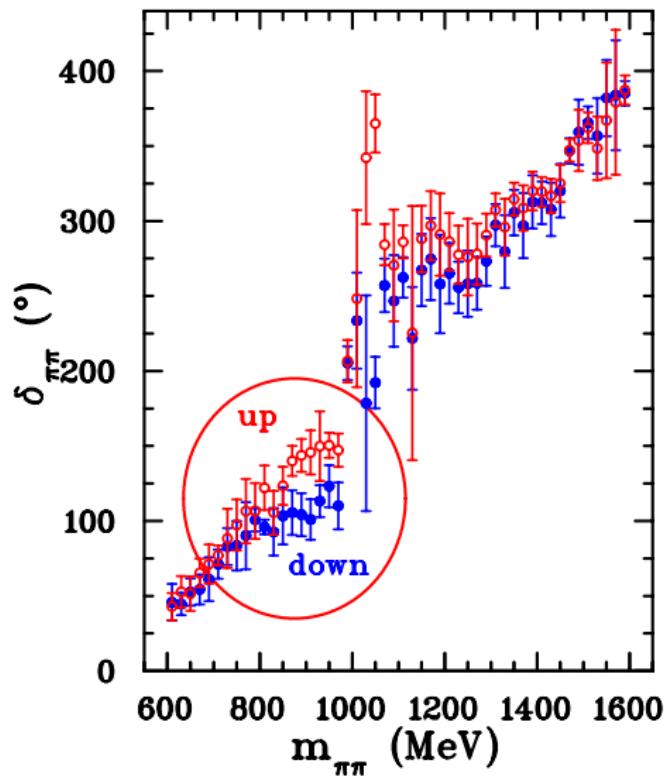
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Condition for crossing symmetry:  $\underline{\text{Re } f_\ell^{I(\text{OUT})}(s) - \text{Re } f_\ell^{I(\text{IN})}(s) = 0}$

# Experimental data on the $\pi\pi$ interactions in the S0 wave (notation $J1$ )

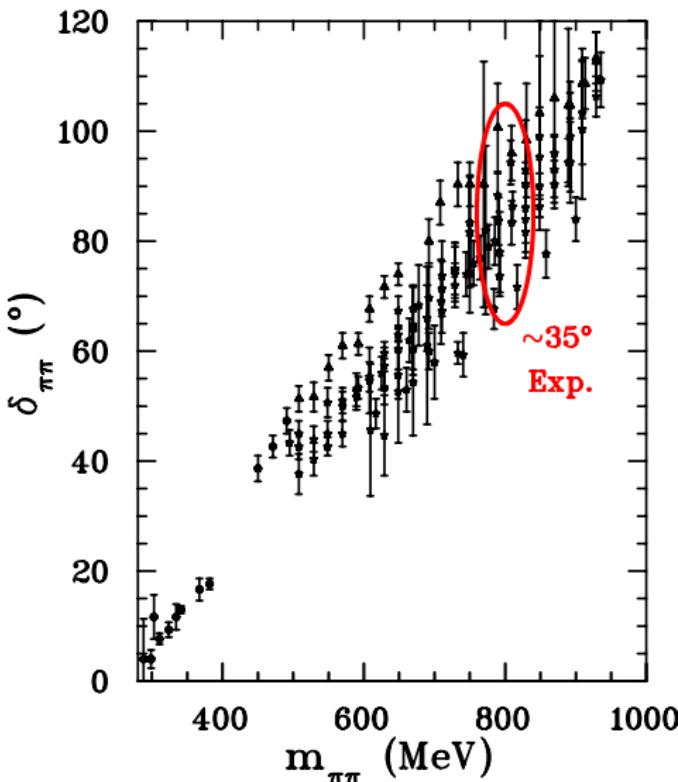


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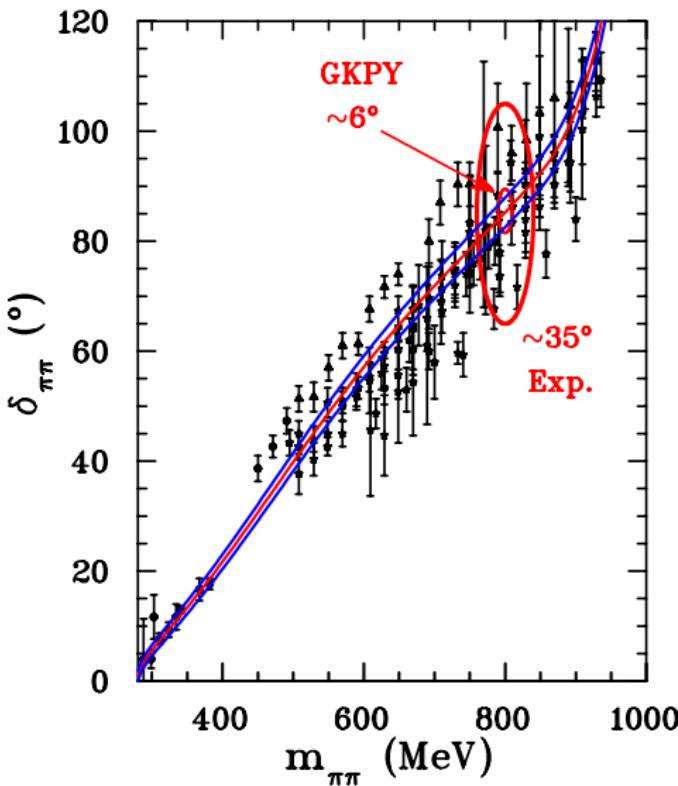
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- "up-down" ambiguity eliminated (in 2003) using the Roy equations (Roy 1971) (two subtractions)
- once subtracted dispersion relations GKY (presented in 2011) - much smaller errors of the output amplitude
- elimination of several sets of experimental data

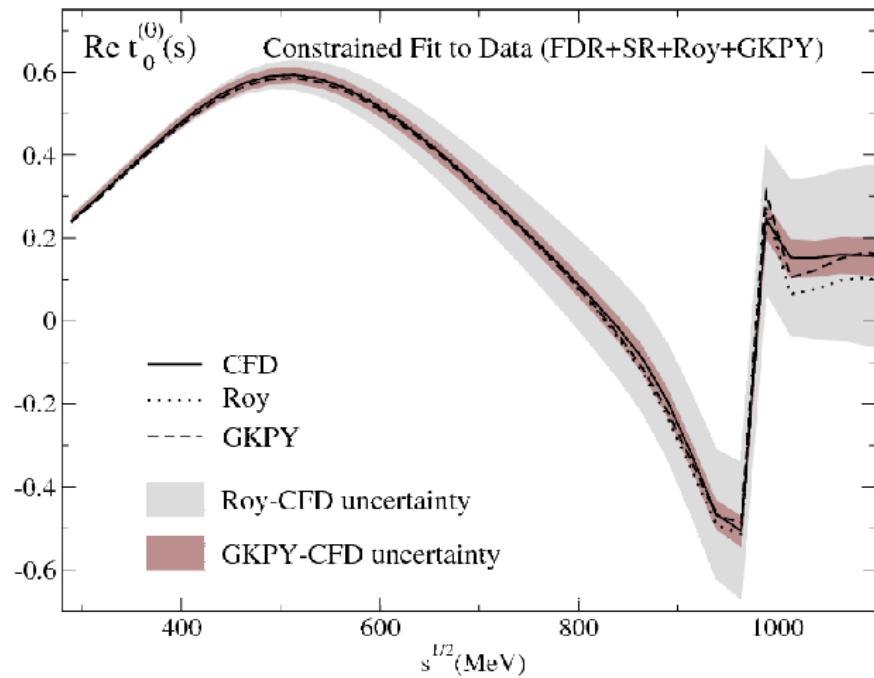


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# precision of the Roy and GKY equations



*"Precise determination of the  $f_0(600)$  and  $f_0(980)$  pole parameters from a dispersive data analysis",*

R. Garcia-Martin, R. Kamiński, J.R. Peláez, J. Ruiz de Elvira,  
Phys. Rev. Lett. 107 (2011) 072001

# Method of combined analysis (data dispersion relations)

- input amplitudes for the  $S, P, D$  and  $F$  waves constructed, at the beginning, only by fit to the data,
- just simple polynomials in energy<sup>2</sup>,
- no assumption on the low threshold parameters (we use NA48/2 data),
- set of dispersion relations used in the analysis:
  - *once subtracted dispersion relations (GKPY)*,
  - *twice subtracted dispersion relations (Roy)*,
  - *Forward Dispersion Relations (FDR)*,
  - *Olsson sum rule (SR)*,
- phenomenological input partial amplitudes used up to 1420 MeV,
- above 1420 MeV - Regge parameterizations

# Method of combined analysis (data dispersion relations)

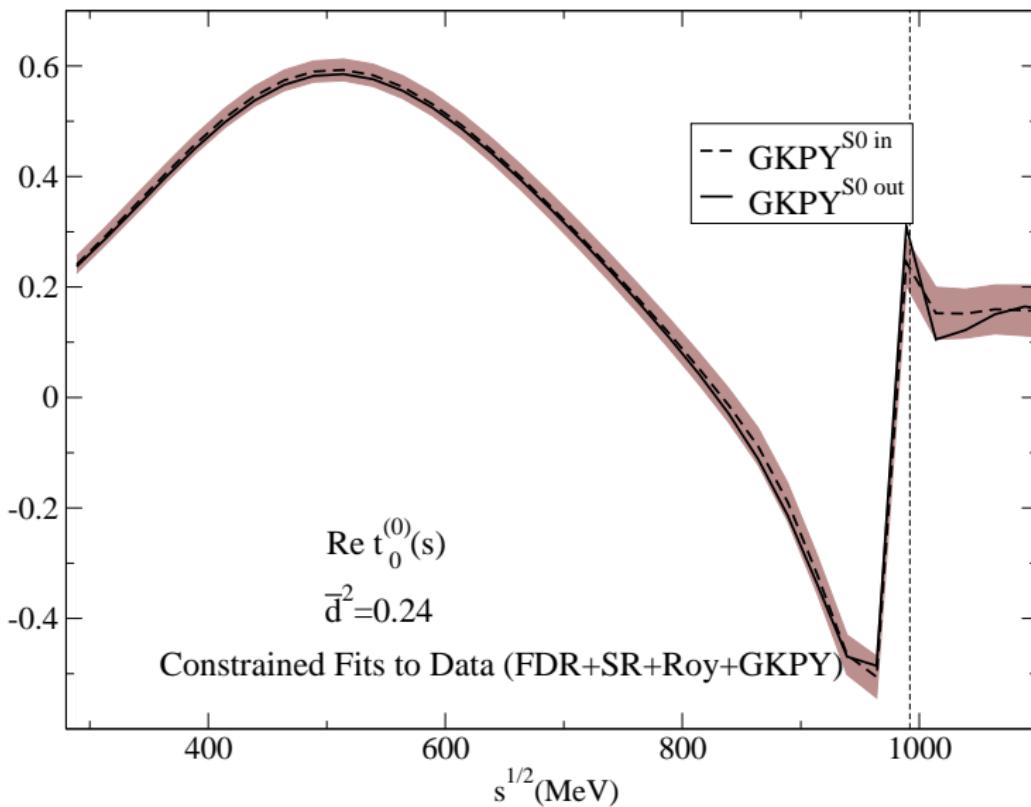
$$\chi^2_{tot} = \chi^2_{data} + \bar{d}_{Roy}^2 + \bar{d}_{GKPY}^2 + \bar{d}_{FDR}^2 + \bar{d}_{SR}^2$$

where  $\bar{d}_i^2 = \frac{1}{\text{number of points}} \sum_j^n \left( \frac{\Delta_i(s_j)}{\delta \Delta_i(s_j)} \right)^2$

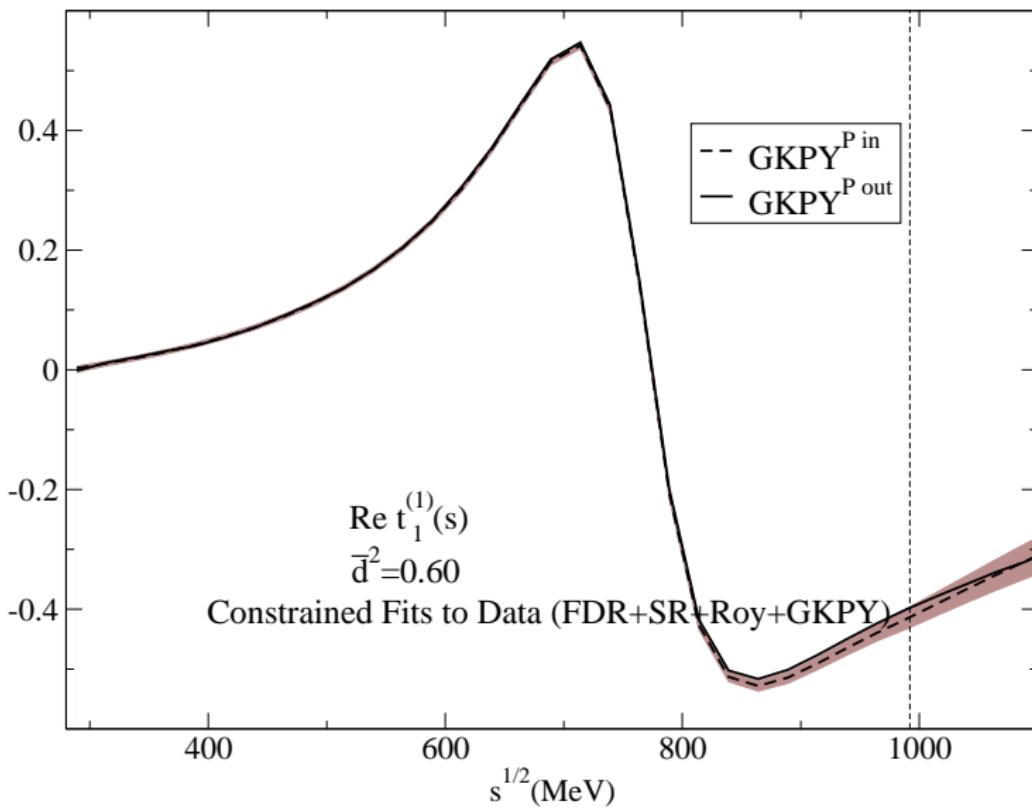
(number of points = 28)

	fit only to data	fit to data and to dispersion relations
$\bar{d}_{Roy}^2$	0.87	0.14
$\bar{d}_{GKPY}^2$	1.9	0.32
$\bar{d}_{FDR}^2$	1.98	0.4

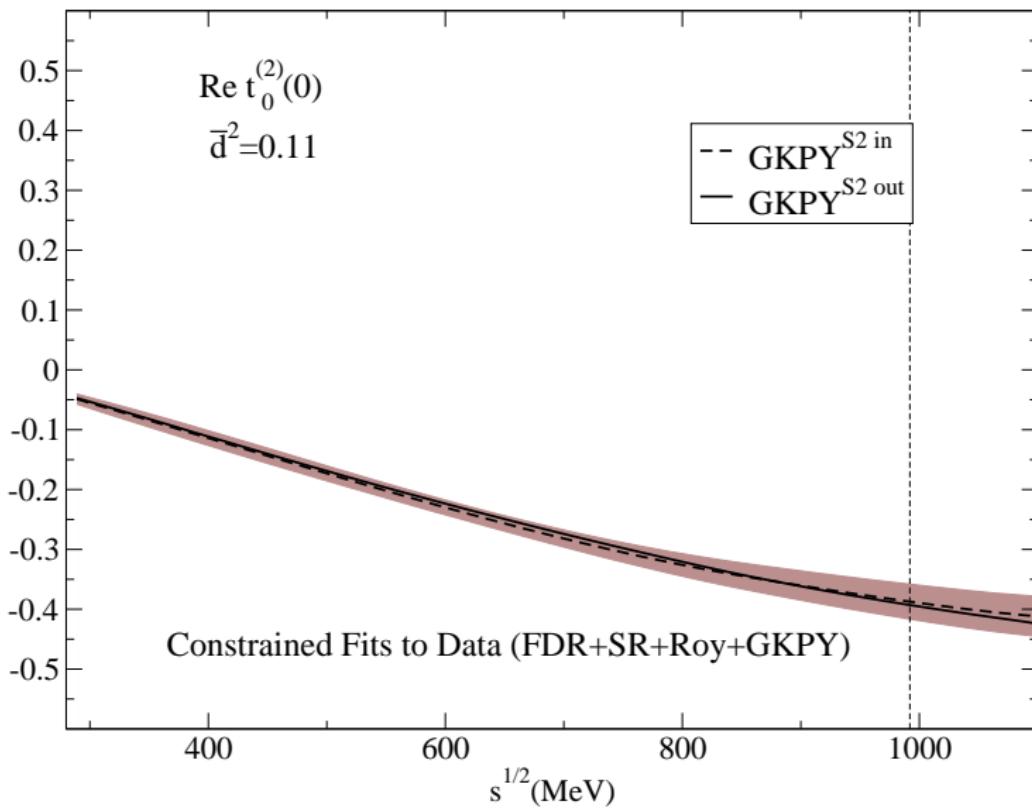
# fit to GKY (S0 wave)



# fit to GKY (P1 wave)



# fit to GKY (S2 wave)



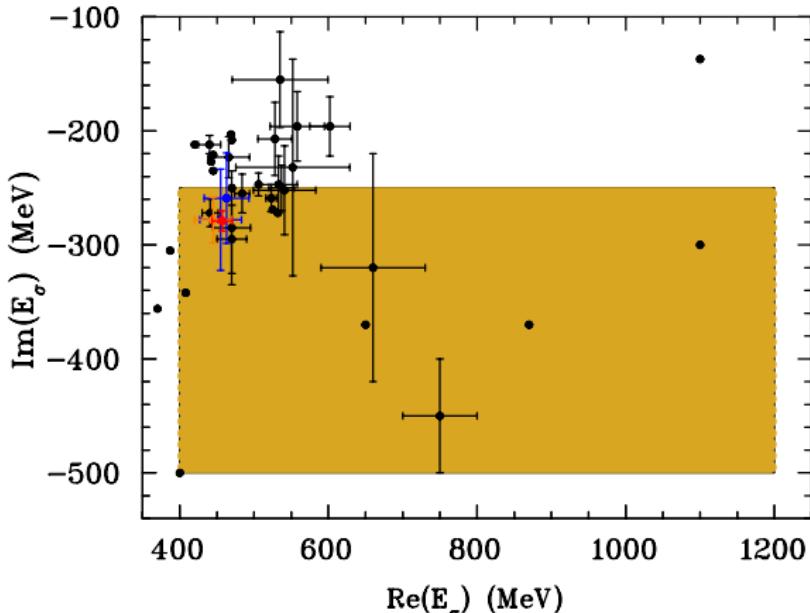
# precise determination of $f_0(600)$ ( $\sigma$ ) meson and threshold parameters

$f_0(600)$  ( $\sigma$ )

- PDG 2010:  
 $M = 400 - 1200$  MeV  
 $\Gamma = 2 \times (250 - 500)$  MeV
- GKPY:  
 $E_\sigma = 457 \pm 14 - i279_{-7}^{+11}$  MeV

threshold parameters, e.g.  $a_0^0$ :

- ChPT + Roy eqs (Bern group):  
 $0.220 \pm 0.005 m_\pi^{-1}$
- GKPY:  
 $0.220 \pm 0.008 m_\pi^{-1}$



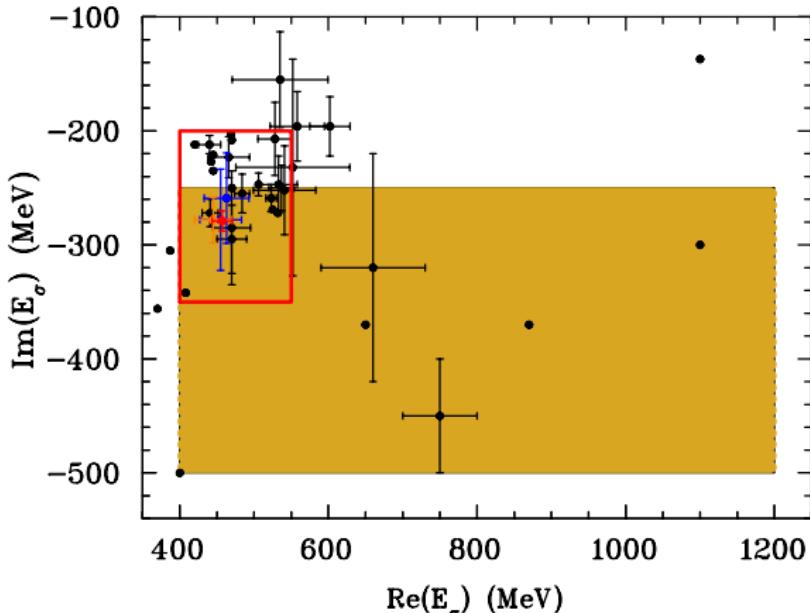
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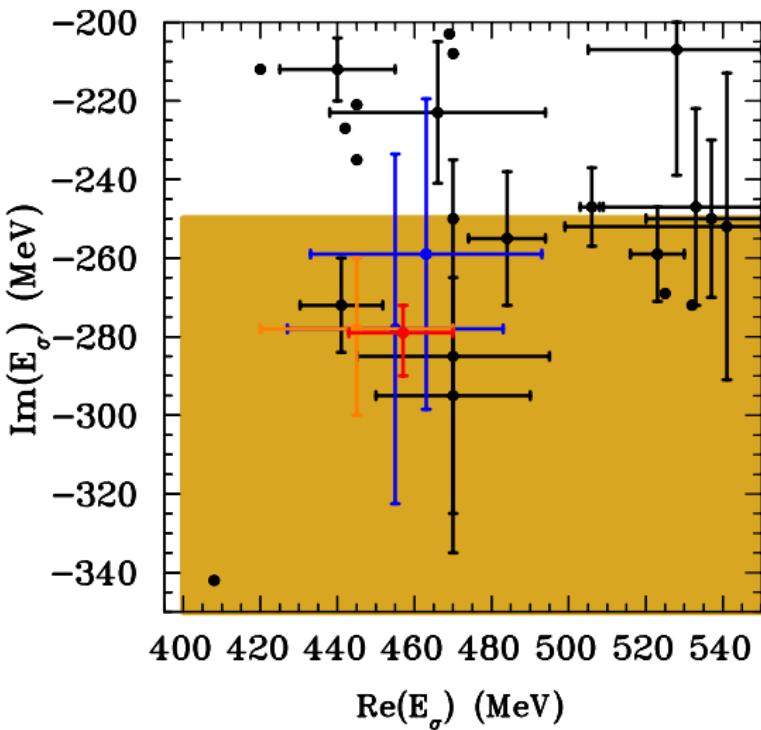
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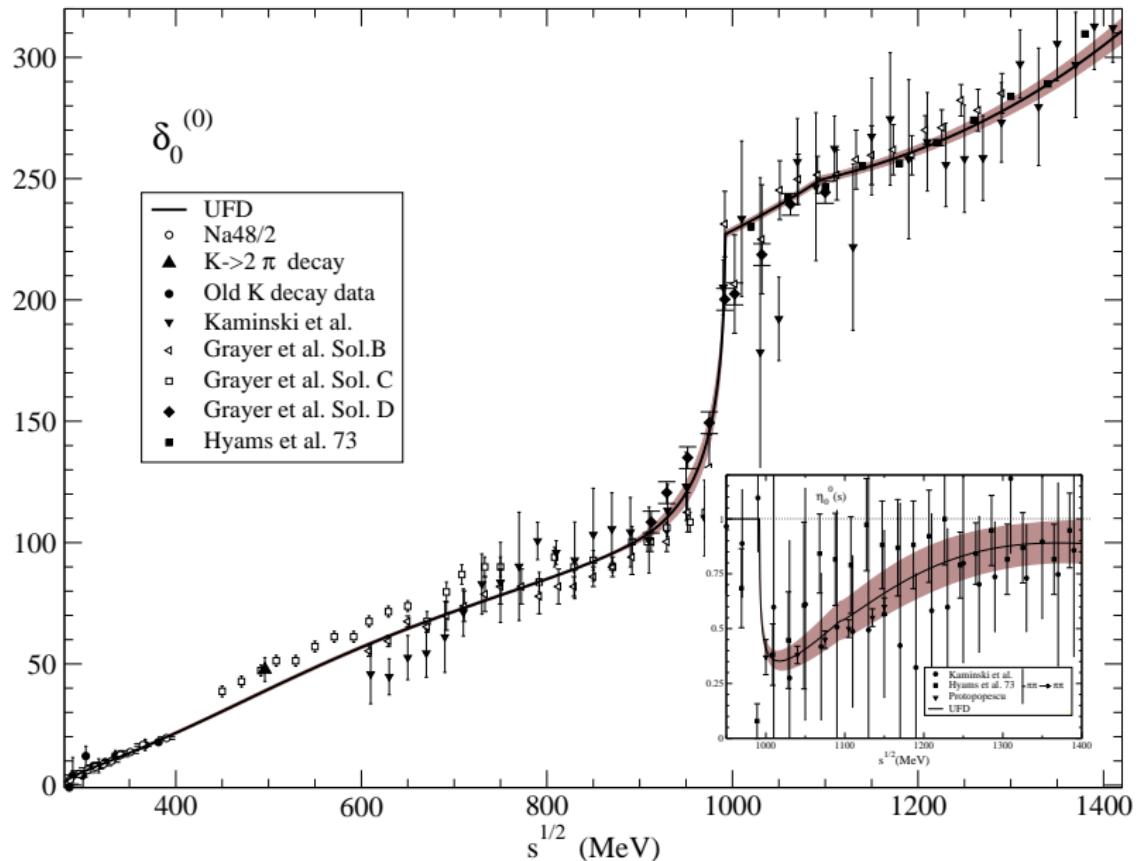
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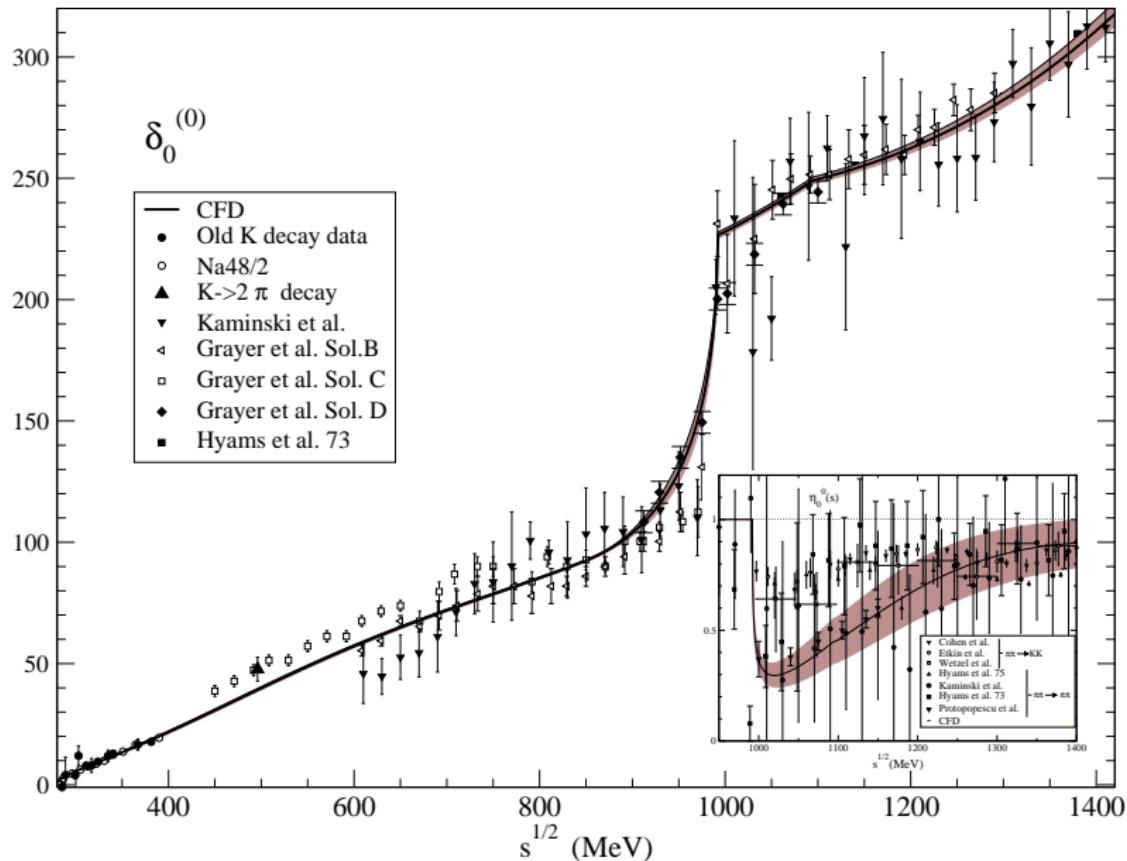
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# precise determination of $f_0(980)$ meson



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Model-independent analytic continuation to the complex plane gets pole at  $s_{pole}^{1/2}$  on the 2nd Riemann sheet

assuming that  $M = Re(s_{pole}^{1/2})$  and  $\Gamma = -2 Im(s_{pole}^{1/2})$  we get:

$$M_{f_0(980)} = 996 \pm 7 \text{ MeV} \text{ and}$$

$$\Gamma_{f_0(980)} = 50^{+20}_{-12} \text{ MeV}$$

PDG'2010: Mass  $m = 980 \pm 10 \text{ MeV}$

Width  $\Gamma = 40 - 100 \text{ MeV}$

# precise determination of couplings to the $\pi\pi$ channel

$$g^2 = -16\pi \lim_{s \rightarrow s_{pole}} (s - s_{pole}) t_\ell(s) (2\ell + 1) / (2p)^{2\ell}$$

where  $p^2 = s/4 - m_\pi^2$ .

	$\sqrt{s_{pole}}$ (MeV)	$ g $
$f_0(600)^{\text{GKPY}}$	$(457^{+14}_{-13}) - i(279^{+11}_{-7})$	$3.59^{+0.11}_{-0.13}$ GeV
$f_0(600)^{\text{Roy}}$	$(445 \pm 25) - i(278^{+22}_{-18})$	$3.4 \pm 0.5$ GeV
$f_0(980)^{\text{GKPY}}$	$(996 \pm 7) - i(25^{+10}_{-6})$	$2.3 \pm 0.2$ GeV
$f_0(980)^{\text{Roy}}$	$(1003^{+5}_{-27}) - i(21^{+10}_{-8})$	$2.5^{+0.2}_{-0.6}$ GeV
$\rho(770)^{\text{GKPY}}$	$(763.7^{+1.7}_{-1.5}) - i(73.2^{+1.0}_{-1.1})$	$6.01^{+0.04}_{-0.07}$
$\rho(770)^{\text{Roy}}$	$(761^{+4}_{-3}) - i(71.7^{+1.9}_{-2.3})$	$5.95^{+0.12}_{-0.08}$

# Conclusions

- due to works on once and twice subtracted dispersion relations with imposed crossing symmetry condition we have in disposal very efficient set of rules **for testing** the partial  $\pi\pi$  amplitudes in the  $S, P, D$  and  $F$  waves,
- we also have set of model independent unitary  $\pi\pi$  amplitudes in those waves in the range from  $2m_\pi$  to several GeV fulfilling very well crossing symmetry below  $\sim 1100$  MeV,
- as an artefact we got very precise values of parameters for the  $f_0(600)$  and  $f_0(980)$  resonances